

Cosmology

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25 March 2014

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 - verbal averaging: can we do better?

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- scalar (GR) averaging: statistically homogeneous spatial slices

verbal averaging

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1. assume homogeneity and isotropy
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verbal averaging

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- practical meaning:

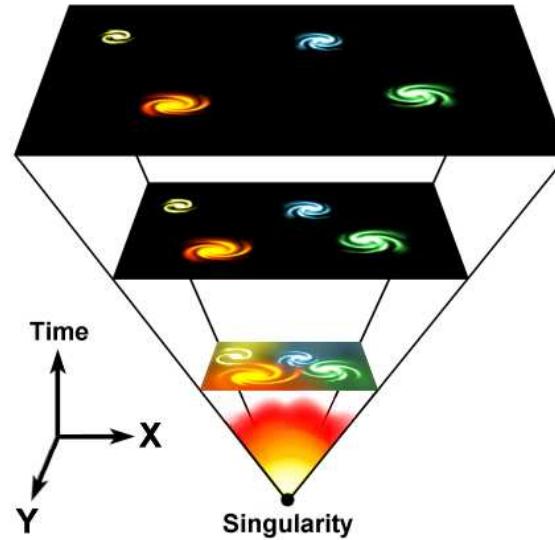
1. assume homogeneity and isotropy
2. find the (differential 4-pseudo-manifold, metric) pairs (M, g) that solve $\mathbf{G} = 8\pi\mathbf{T}$
3. assume that (M, g) remains unchanged if we add density perturbations to an early time slice

verbal averaging

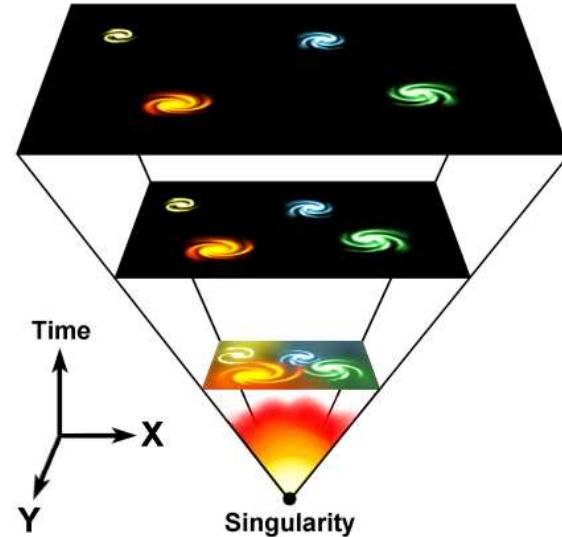
- w:Comoving coordinates

verbal averaging

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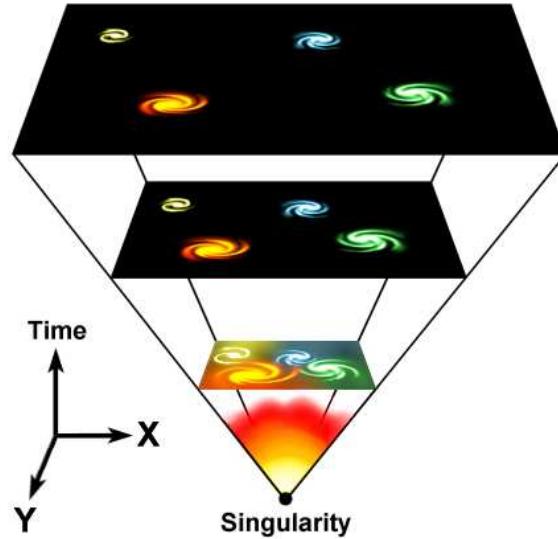
verbal averaging



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$$\Delta x(t) = a(t)\Delta r$$

verbal averaging



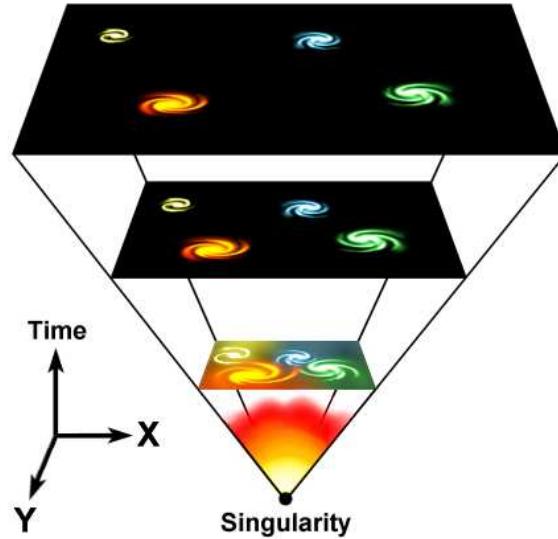
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- spherical coordinates for spatial slice

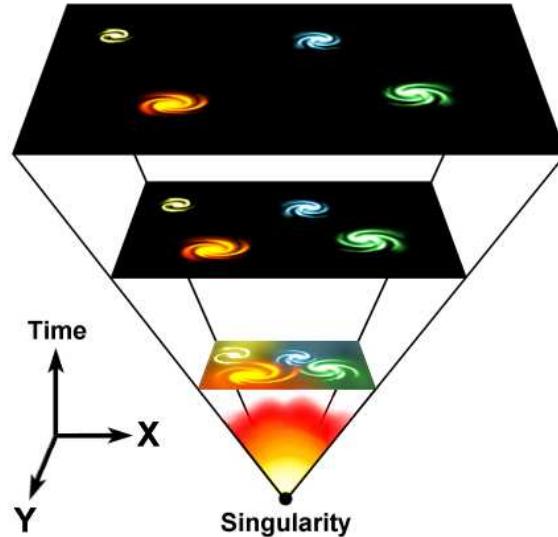
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$$\int_{(t,r_1,\theta,\phi)}^{(t,r_2,\theta,\phi)} ds = a(t)\Delta r = a(t)|r_2 - r_1|$$

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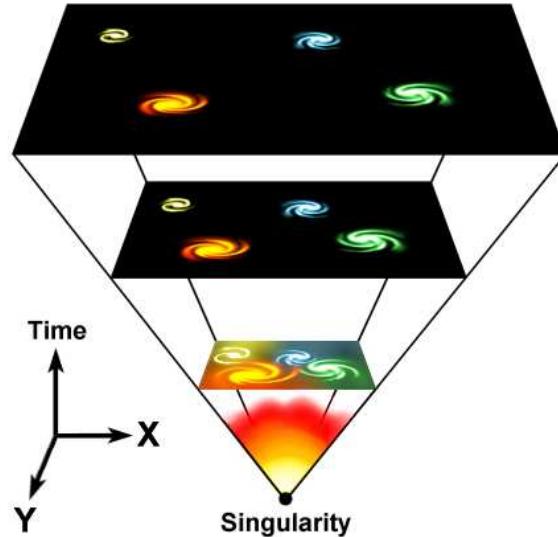


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where all expansion/contraction \rightarrow w:scale factor $a(t)$

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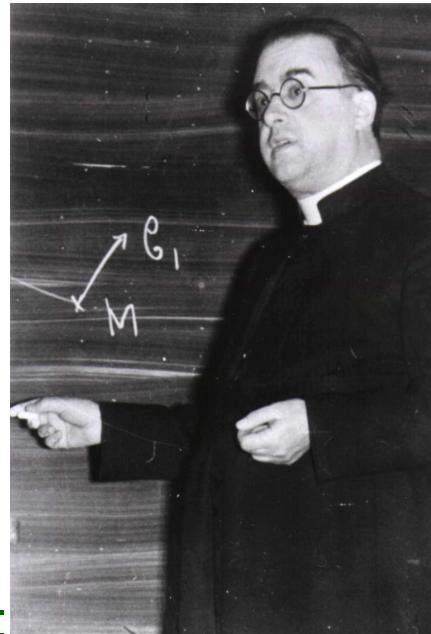
- universe is static in comoving coordinates (r, θ, ϕ)

FLRW metric

- [w:Friedmann–Lemaître–Robertson–Walker metric](#)

FLRW metric

■ w:Friedmann–Lemaître–Robertson–Walker metric



- w: Alexander Friedmann w: Howard Percy Robertson
w: Arthur Geoffrey Walker

FLRW metric

$$ds^2 = -dt^2 + \dots$$

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$$u := \int \frac{dt}{a(t)}$$

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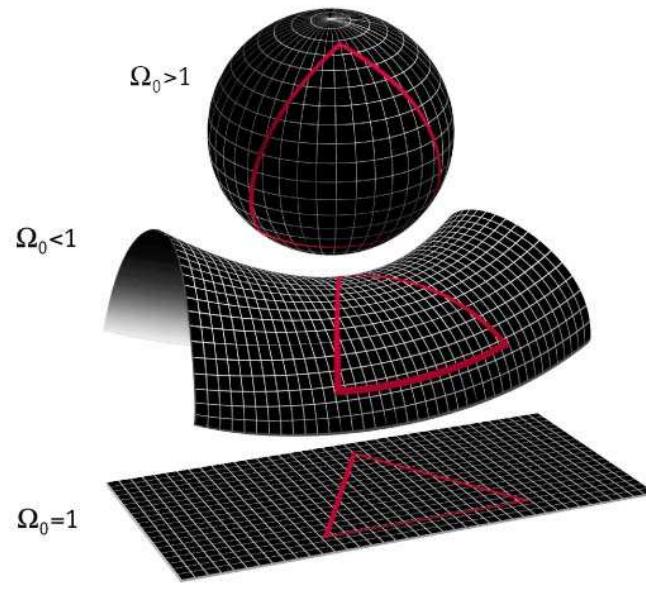
\Rightarrow

$$ds^2 = a^2(u) [-du^2 + dr^2 + r_\perp^2(d\theta^2 + \cos^2\theta d\phi^2)]$$

- but $\int du \neq$ proper time; *more:* [arXiv:astro-ph/0707.2106](https://arxiv.org/abs/astro-ph/0707.2106)

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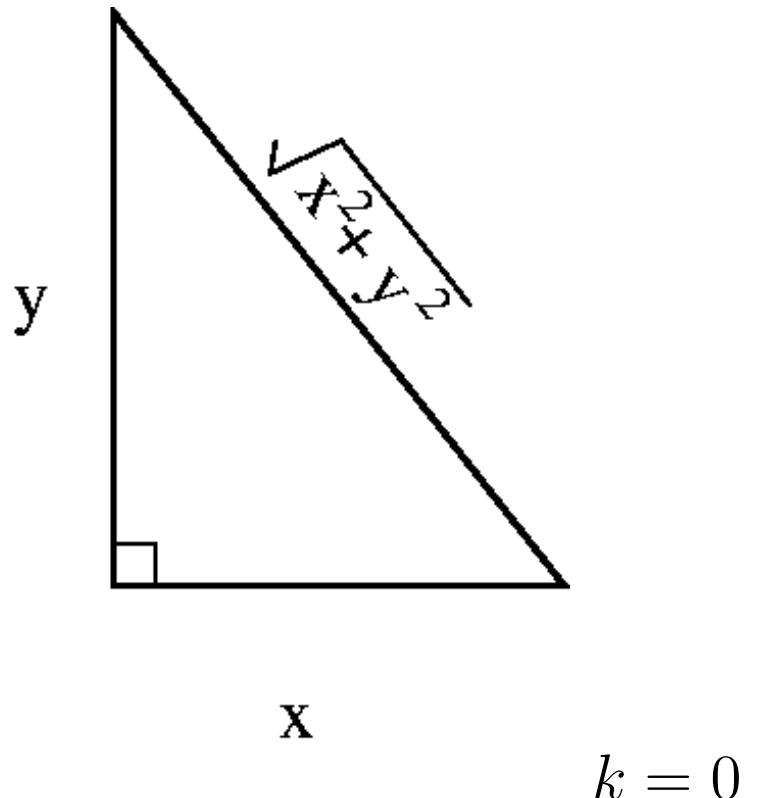


where $r_\perp := \begin{cases} R_C \sinh \frac{r}{R_C} & k < 0 \\ r & k = 0 \\ R_C \sin \frac{r}{R_C} & k > 0 \end{cases}$

for a comoving radius of curvature R_C and curvature of sign k

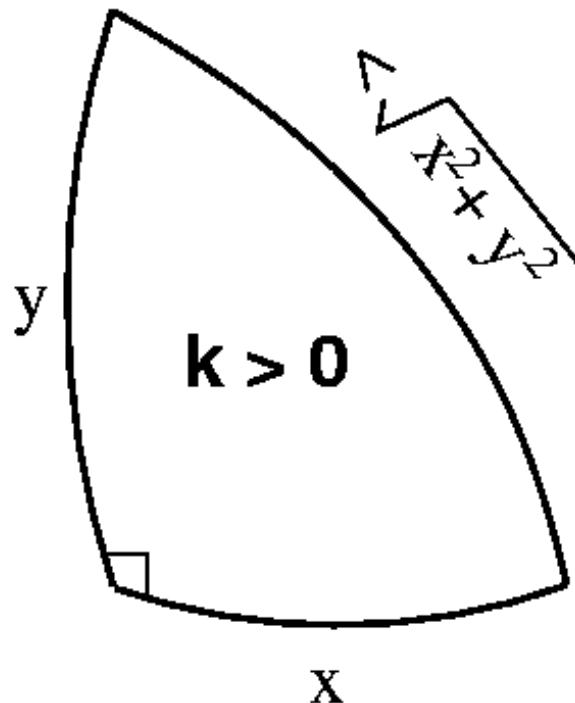
curvature

- on a spatial slice (fixed value of t):



curvature

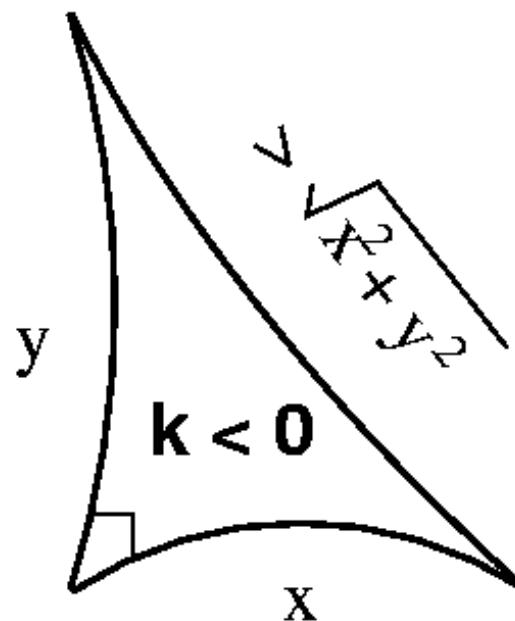
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$$k > 0$$

curvature

- on a spatial slice (fixed value of t):



$$k < 0$$

2D curvature intuition: $k > 0$

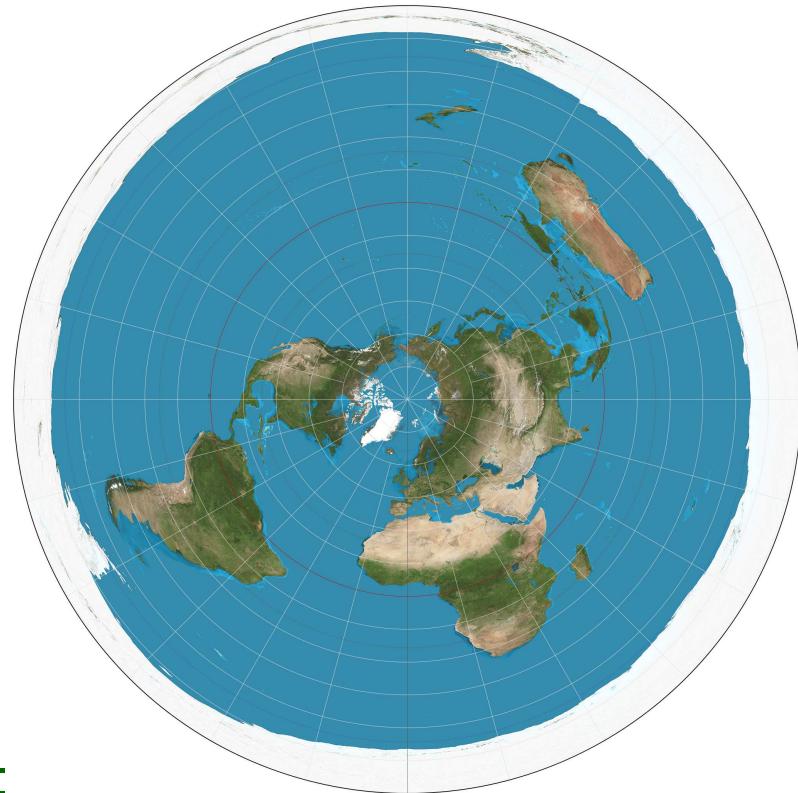
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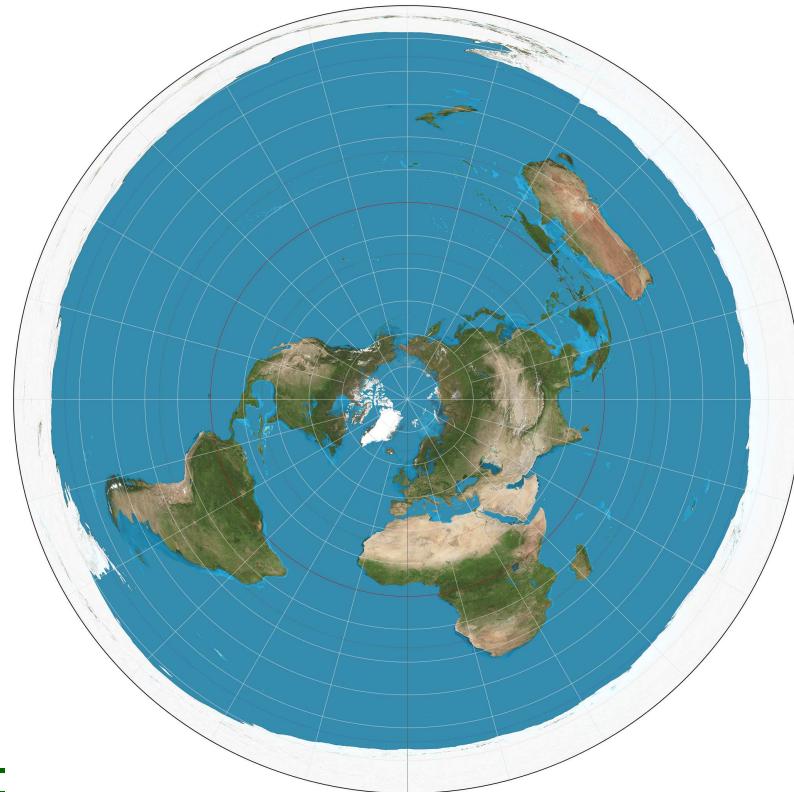


w:

(al-Biruni, c. 1000 CE)

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- intuition switch: S^2 easier vs S^3 more physical

2D topology intuition ($k = 0$)



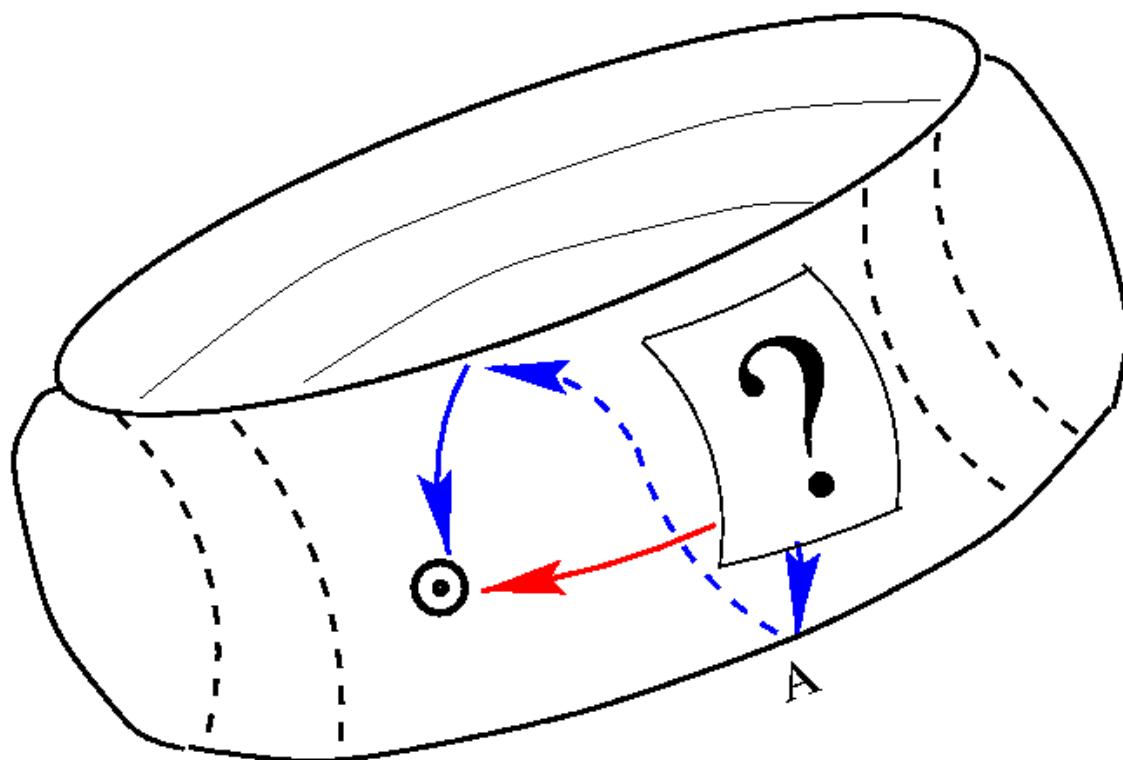
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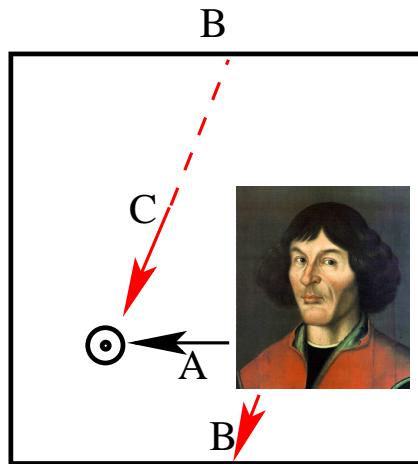


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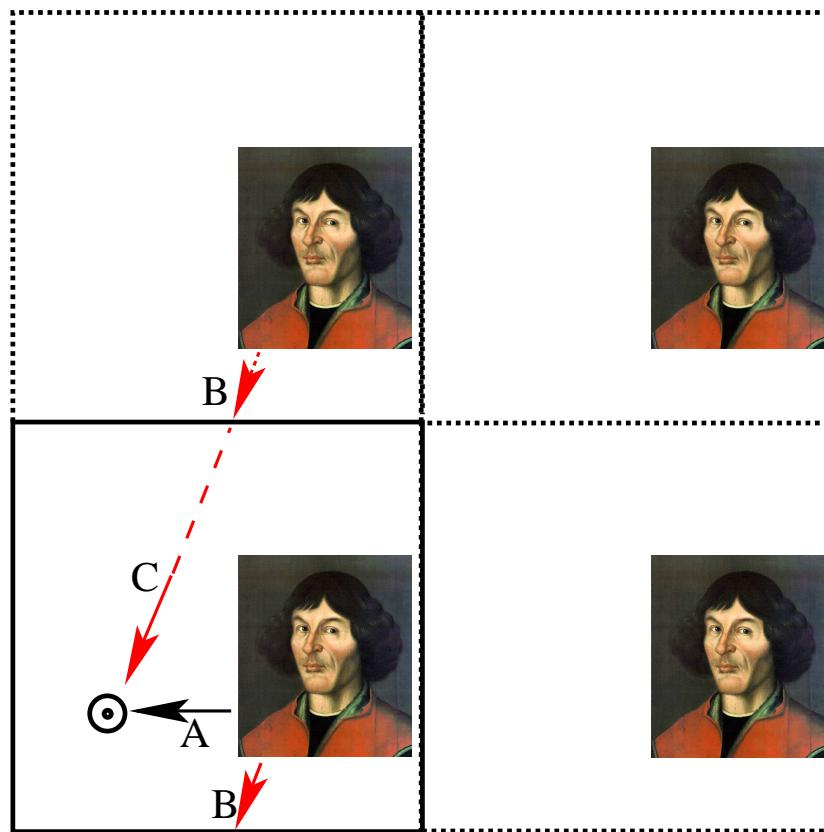


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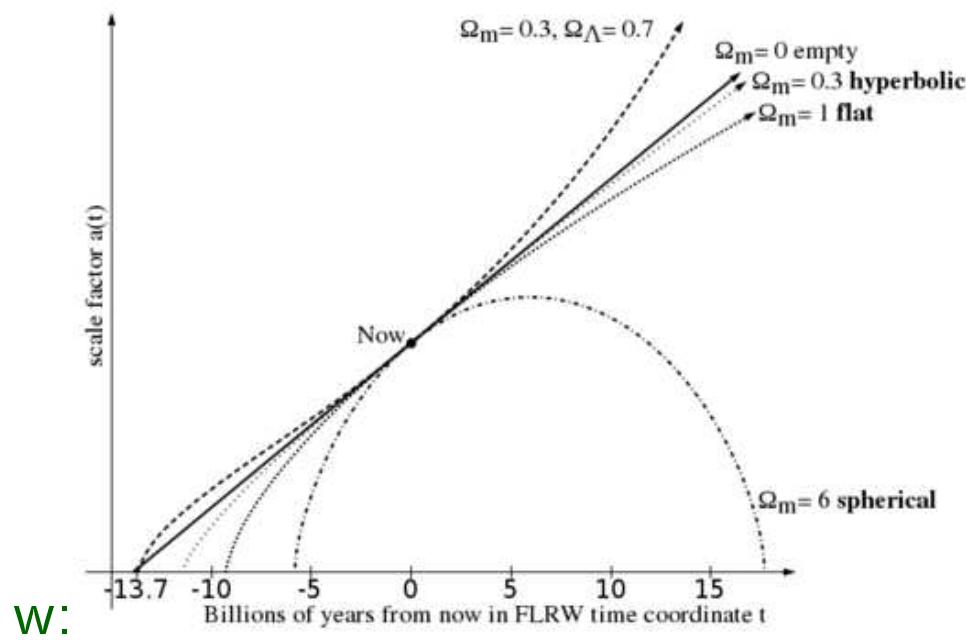
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(Defn: $a_0 := 1$)

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$$1 + z = \frac{1}{a_{\text{em}}}$$

(Defn of redshift z)

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- radiation density: $E = h\nu \Rightarrow \rho_r \propto a^{-4} = (1 + z)^4$

Black body: COBE (~ 1992)

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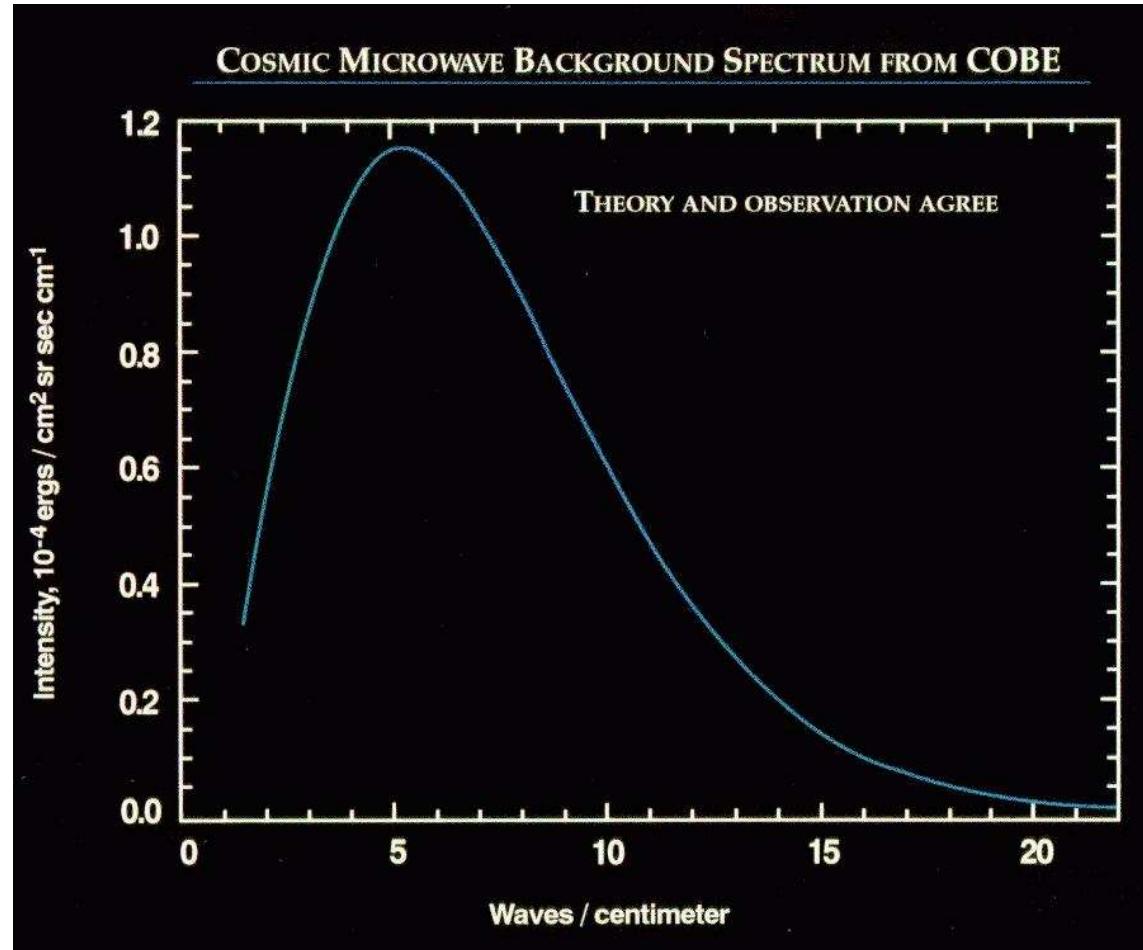
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Black body: COBE (~ 1992)

- COBE /FIRAS (Far Infrared Absolute Spectrophotometer)

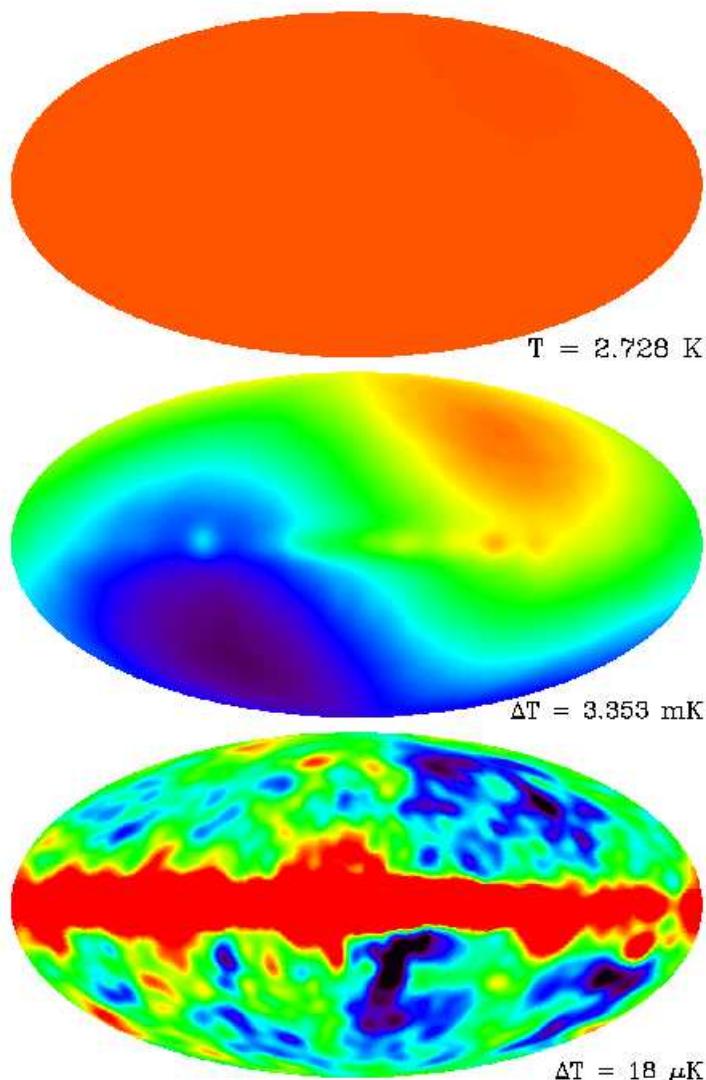
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- COBE /FIRAS (Far Infrared Absolute Spectrophotometer)



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- COBE /DMR (Differential Microwave Radiometer)

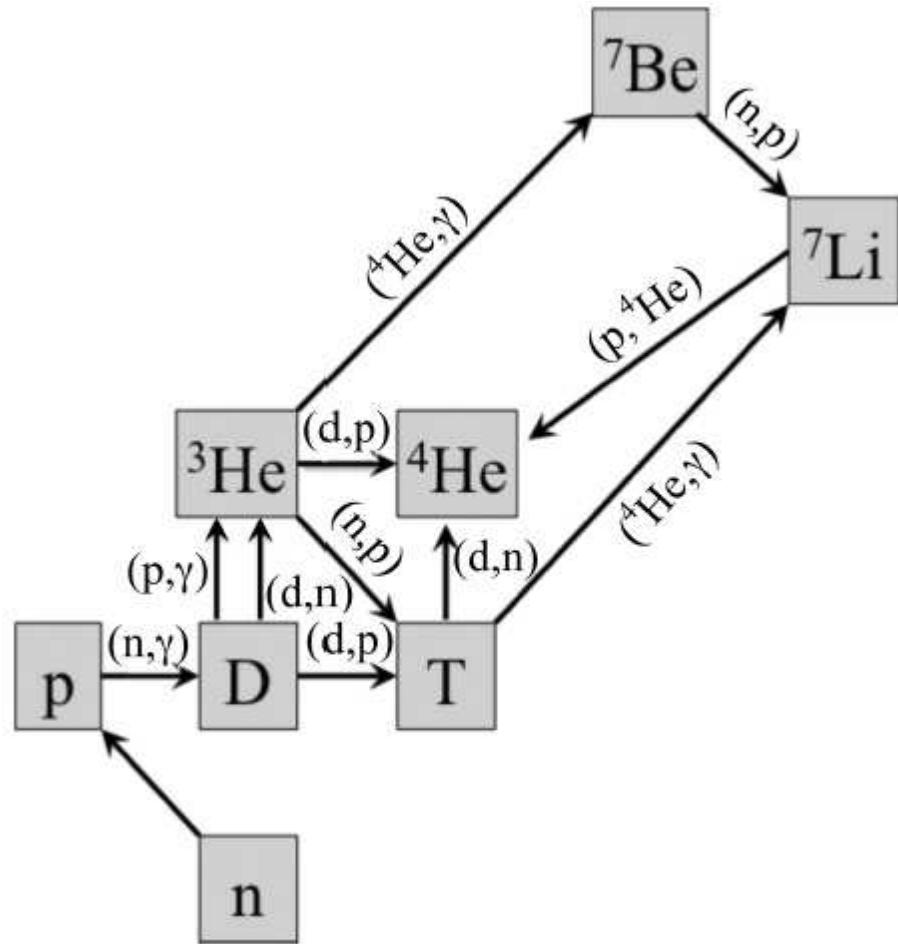


BBN: Big bang nucleosynthesis

■ Alpher, Bethe, & Gamow (1948; ADS:1948PhRv...73..803A)

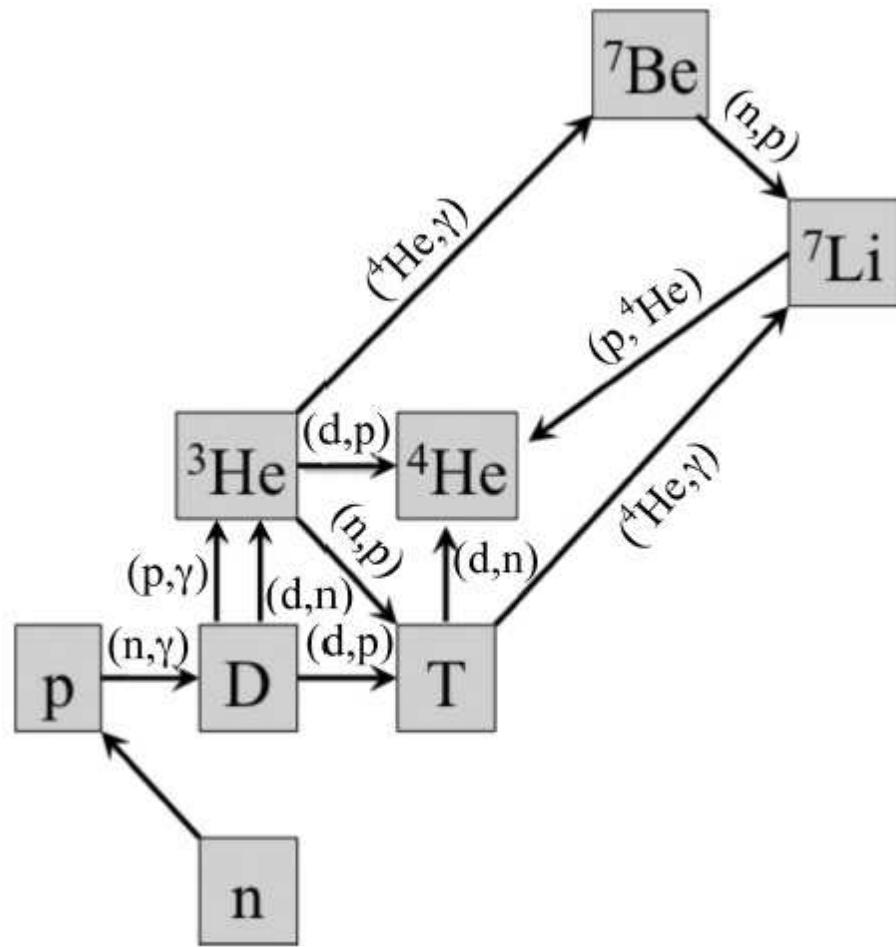
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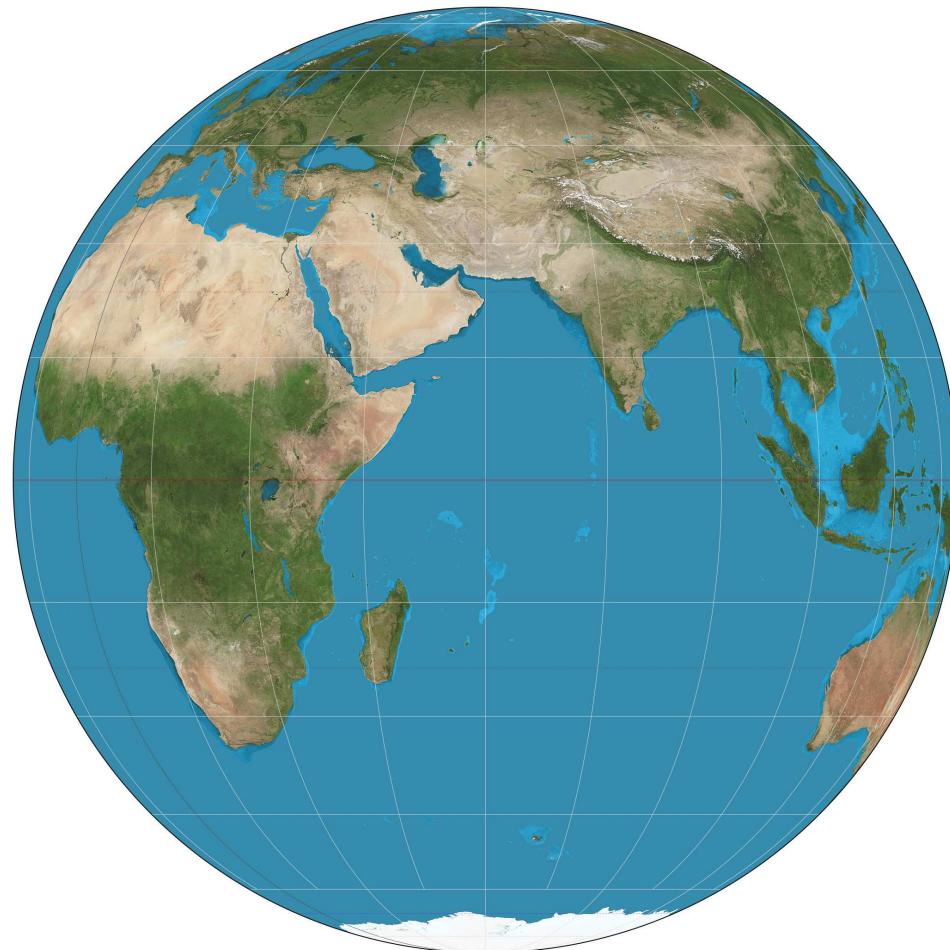
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Friedmann Eqn:

$$\frac{c^2 k}{a^2} + \frac{\dot{a}^2}{a^2} = \frac{8\pi G \rho}{3}$$

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■

acceleration Eqn:

$$\frac{\ddot{a}}{a} = -\frac{4\pi G (\rho + 3p/c^2)}{3}$$

FLRW matter-dominated epoch

■ Friedmann Eqn:

$$\frac{\dot{a}^2}{a^2} = \frac{8\pi G \rho}{3} - \frac{c^2 k}{a^2}$$

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FLRW matter-dominated epoch

- Friedmann Eqn:

$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch: $\rho = \rho_m = \rho_{m0} a^{-3}$
- $k = 0$ case: $\dot{a} \propto a^{-1/2}$ for $a > 0$

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- $k = 0$ case: $da \propto a^{-1/2} dt$ for $a > 0$

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$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$
- matter-dominated epoch: $\rho = \rho_m = \rho_{m0} a^{-3}$
- $k = 0$ case: $a^{1/2} da \propto dt$

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- matter-dominated epoch: $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$ case: $(2/3)a^{3/2} \propto t$

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- Friedmann Eqn:
$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$
- matter-dominated epoch: $\rho = \rho_m = \rho_{m0} a^{-3}$
- $k = 0$ case: $a \propto t^{2/3}$

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$$\dot{a}^2 = \frac{8\pi G \rho_{m0}}{3a} - c^2 k$$

- matter-dominated epoch: $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$ case:

$$a = \left(\frac{t}{t_0}\right)^{2/3}$$

Einstein–de Sitter model (EdS)

FLRW matter-dominated epoch

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Einstein–de Sitter model (EdS)

Defn: Hubble parameter $H := \dot{a}/a$

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Defn: Hubble parameter $H := \dot{a}/a$

\Rightarrow Hubble constant $H_0 := H(z = 0)$

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- $k = 0$ case: $a = \left(\frac{t}{t_0}\right)^{2/3}$ Einstein–de Sitter model (EdS)

Defn: Hubble parameter $H := \dot{a}/a$

\Rightarrow Hubble constant $H_0 := H(z = 0)$

\Rightarrow $k = 0$ case: $\frac{\dot{a}}{a} = \frac{2}{3t}$

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$\Rightarrow H(t) = \frac{2}{3t};$

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Defn: Hubble parameter $H := \dot{a}/a$

\Rightarrow Hubble constant $H_0 := H(z=0)$

\Rightarrow $k = 0$ case: $\frac{\dot{a}}{a} = \frac{2}{3t}$

$\Rightarrow H(t) = \frac{2}{3t}; H_0 = H(t_0) = \frac{2}{3t_0}$

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- matter-dominated epoch: $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$ case: $a = \left(\frac{t}{t_0}\right)^{2/3}$ Einstein–de Sitter model (EdS)

$$\Rightarrow H(t) = \frac{2}{3t}; H_0 = H(t_0) = \frac{2}{3t_0}$$

- convenient conversion: $1 \text{ km/s} \approx 1.04 \text{ kpc/Gyr} \approx 1 \text{ kpc/Gyr}$

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$$\Rightarrow H(t) = \frac{2}{3t}; H_0 = H(t_0) = \frac{2}{3t_0}$$

- Lemaître (1927) ADS:1927ASSB...47...49L: $H_0 \approx 600 \text{ km/s/Mpc}$

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- Lemaître (1927) ADS:1927ASSB...47...49L: $H_0 \approx 0.6 \text{ Gyr}^{-1}$

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- Lemaître (1927) ADS:1927ASSB...47...49L: $H_0 \approx 0.6 \text{ Gyr}^{-1}$
- Hubble (1929) ADS:1929PNAS...15..168H: $H_0 \approx 500 \text{ km/s/Mpc}$

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$$\Rightarrow \text{EdS would give } t_0 = \frac{2}{3H_0}$$

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- matter-dominated epoch: $\rho = \rho_m = \rho_{m0} a^{-3}$

- $k = 0$ case: $a = \left(\frac{t}{t_0}\right)^{2/3}$ Einstein–de Sitter model (EdS)

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- Lemaître (1927) ADS:1927ASSB...47...49L: $H_0 \approx 0.6 \text{ Gyr}^{-1}$

- Hubble (1929) ADS:1929PNAS...15..168H: $H_0 \approx 0.5 \text{ Gyr}^{-1}$

$$\Rightarrow \text{EdS would give } t_0 = \frac{2}{3H_0} \approx 1.3 \text{ Gyr}$$

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$$\Rightarrow \text{EdS would give } t_0 = \frac{2}{3H_0} \approx 1.3 \text{ Gyr} < t_{\text{Earth}} \approx 4.5 \text{ Gyr}$$

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- 1980's: $H_0 \approx 50$ or 100 km/s/Mpc

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$$\Rightarrow \text{EdS would give } t_0 = \frac{2}{3H_0} \approx 1.3 \text{ Gyr} < t_{\text{Earth}} \approx 4.5 \text{ Gyr}$$

- 1980's: $H_0 \approx 0.05$ or 0.1 Gyr^{-1}

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- $k = 0$ case: $a = \left(\frac{t}{t_0}\right)^{2/3}$ Einstein–de Sitter model (EdS)

$$\Rightarrow H(t) = \frac{2}{3t}; H_0 = H(t_0) = \frac{2}{3t_0}$$

- Lemaître (1927) ADS:1927ASSB...47...49L: $H_0 \approx 0.6 \text{ Gyr}^{-1}$

- Hubble (1929) ADS:1929PNAS...15..168H: $H_0 \approx 0.5 \text{ Gyr}^{-1}$

$$\Rightarrow \text{EdS would give } t_0 = \frac{2}{3H_0} \approx 1.3 \text{ Gyr} < t_{\text{Earth}} \approx 4.5 \text{ Gyr}$$

- 1980's: $H_0 \approx 0.05$ or $0.1 \text{ Gyr}^{-1} \Rightarrow t_0(\text{EdS}) \approx 13.0$ or 6.5 Gyr , resp.

FLRW: ρ_{crit}

■ Friedmann Eqn:

$$\frac{\dot{a}^2}{a^2} = \frac{8\pi G \rho}{3} - \frac{c^2 k}{a^2}$$

FLRW: ρ_{crit}

■ Friedmann Eqn:

$$H^2 = \frac{8\pi G \rho}{3} - \frac{c^2 k}{a^2}$$

FLRW: ρ_{crit}

- Friedmann Eqn:
$$H^2 = \frac{8\pi G \rho}{3} - \frac{c^2 k}{a^2}$$
- consider a fixed observation, e.g. $H_0 = 100 \text{ km/s/Mpc}$

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- consider a fixed observation, e.g. $H_0 = 100 \text{ km/s/Mpc}$
 - ◆ $\rho_{m0} = \frac{3H_0^2}{8\pi G} \Leftrightarrow k = 0$

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$$H^2 = \frac{8\pi G \rho}{3} - \frac{c^2 k}{a^2}$$
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 - ◆ $\rho_{m0} = \frac{3H_0^2}{8\pi G} \Leftrightarrow k = 0$ flat

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$$H^2 = \frac{8\pi G \rho}{3} - \frac{c^2 k}{a^2}$$
- consider a fixed observation, e.g. $H_0 = 100 \text{ km/s/Mpc}$
 - ◆ $\rho_{m0} = \frac{3H_0^2}{8\pi G} \Leftrightarrow k = 0$ flat
 - ◆ $\rho_{m0} > \frac{3H_0^2}{8\pi G} \Leftrightarrow k > 0$

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 - ◆ $\rho_{m0} = \frac{3H_0^2}{8\pi G} \Leftrightarrow k = 0$ flat
 - ◆ $\rho_{m0} > \frac{3H_0^2}{8\pi G} \Leftrightarrow k > 0$ spherical

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 - ◆ $\rho_{m0} < \frac{3H_0^2}{8\pi G} \Leftrightarrow k < 0$

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- Friedmann Eqn:
$$H^2 = \frac{8\pi G \rho}{3} - \frac{c^2 k}{a^2}$$
- consider a fixed observation, e.g. $H_0 = 100 \text{ km/s/Mpc}$
 - ◆ $\rho_{m0} = \frac{3H_0^2}{8\pi G} \Leftrightarrow k = 0$ flat
 - ◆ $\rho_{m0} > \frac{3H_0^2}{8\pi G} \Leftrightarrow k > 0$ spherical
 - ◆ $\rho_{m0} < \frac{3H_0^2}{8\pi G} \Leftrightarrow k < 0$ hyperbolic

FLRW: ρ_{crit}

■ Friedmann Eqn:

$$H^2 = \frac{8\pi G \rho}{3} - \frac{c^2 k}{a^2}$$

FLRW: ρ_{crit}

■ Friedmann Eqn:
$$H^2 = \frac{8\pi G \rho}{3} - \frac{c^2 k}{a^2}$$

Defn:
$$\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$$
 critical density

FLRW: ρ_{crit}

■ Friedmann Eqn:

$$H^2 = \frac{\rho H^2}{\rho_{\text{crit}}} - \frac{c^2 k}{a^2}$$

Defn: $\boxed{\rho_{\text{crit}} := \frac{3H^2}{8\pi G}}$ critical density

FLRW: ρ_{crit}

■ Friedmann Eqn:

$$H^2 = \frac{\rho H^2}{\rho_{\text{crit}}} - \frac{c^2 k}{a^2}$$

Defn: $\boxed{\rho_{\text{crit}} := \frac{3H^2}{8\pi G}}$ critical density

Defn: $\boxed{\Omega_m := \frac{\rho}{\rho_{\text{crit}}}}$ matter density parameter

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$$H^2 = \Omega_m H^2 - \frac{c^2 k}{a^2}$$

Defn:
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 critical density

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Defn:
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 critical density

Defn:
$$\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$$
 matter density parameter

Defn:
$$\Omega_k := -\frac{c^2 k}{a^2 H^2}$$
 curvature density parameter

FLRW: ρ_{crit}

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$$H^2 = \Omega_m H^2 - \frac{c^2 k}{a^2}$$

Defn:
$$\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$$
 critical density

Defn:
$$\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$$
 matter density parameter

Defn:
$$\Omega_k := -\frac{c^2 k}{a^2 H^2}$$
 curvature density parameter (sign reversal!)

FLRW: ρ_{crit}

■ Friedmann Eqn: $H^2 = \Omega_m H^2 + \Omega_k H^2$

Defn: $\boxed{\rho_{\text{crit}} := \frac{3H^2}{8\pi G}}$ critical density

Defn: $\boxed{\Omega_m := \frac{\rho}{\rho_{\text{crit}}}}$ matter density parameter

Defn: $\boxed{\Omega_k := -\frac{c^2 k}{a^2 H^2}}$ curvature density parameter (sign reversal!)

FLRW: ρ_{crit}

■ Friedmann Eqn:

$$1 = \Omega_m + \Omega_k$$

Defn: $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$ critical density

Defn: $\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$ matter density parameter

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- consider a fixed observation, e.g. $H_0 = 100$ km/s/Mpc

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Defn: $\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$ matter density parameter

Defn: $\Omega_k := -\frac{c^2 k}{a^2 H^2}$ curvature density parameter (sign reversal!)

- consider a fixed observation, e.g. $H_0 = 100 \text{ km/s/Mpc}$

- ◆ $\Omega_{m0} = 1 \Leftrightarrow k = 0$

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■ consider a fixed observation, e.g. $H_0 = 100 \text{ km/s/Mpc}$

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- ◆ $\Omega_{m0} > 1 \Leftrightarrow k > 0$

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$$1 = \Omega_m + \Omega_k$$

Defn: $\rho_{\text{crit}} := \frac{3H^2}{8\pi G}$ critical density

Defn: $\Omega_m := \frac{\rho}{\rho_{\text{crit}}}$ matter density parameter

Defn: $\Omega_k := -\frac{c^2 k}{a^2 H^2}$ curvature density parameter (sign reversal!)

■ consider a fixed observation, e.g. $H_0 = 100 \text{ km/s/Mpc}$

- ◆ $\Omega_{m0} = 1 \Leftrightarrow k = 0$ flat
- ◆ $\Omega_{m0} > 1 \Leftrightarrow k > 0$ spherical
- ◆ $\Omega_{m0} < 1 \Leftrightarrow k < 0$ hyperbolic

FLRW: ρ_{crit}

■ Friedmann Eqn:

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■ $\Omega_{\text{tot}} :=$

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- $\Omega_{\text{tot}} := \Omega_m +$

FLRW: ρ_{crit}

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■ $\Omega_{\text{tot}} := \Omega_m + \Omega_r +$

FLRW: ρ_{crit}

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■ $\Omega_{\text{tot}} := \Omega_m + \Omega_r + \Omega_\Lambda$

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■ $\Omega_{\text{tot}} := \Omega_b + \Omega_{\text{nbDM}} + \Omega_r + \Omega_\Lambda$

FLRW curvature constant

- metric in
 - ◆ azimuthal equidistant coords: R_C

FLRW curvature constant

■ metric in

- ◆ azimuthal equidistant coords: R_C
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FLRW curvature constant

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■ orthographic: $1 - kr^2 = 0$ coord singularity at equator

FLRW curvature constant

- metric in
 - ◆ azimuthal equidistant coords: R_C
 - ◆ orthographic coords: k
- orthographic: $1 - kr^2 = 0$ coord singularity at equator
- $\Rightarrow kR_C^2 = 1$

FLRW curvature constant

- metric in
 - ◆ azimuthal equidistant coords: R_C
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- $\Rightarrow kR_C^2 = 1 \Rightarrow k = 1/R_C^2$

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■ orthographic: $1 - kr^2 = 0$ coord singularity at equator

■ $\Rightarrow kR_C^2 = 1 \Rightarrow k = 1/R_C^2$

■ Defn: $\Omega_k := -\frac{c^2 k}{a^2 H^2}$ $\Rightarrow \Omega_{k0} = -\frac{c^2 k}{H_0^2}$

FLRW curvature constant

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- ◆ azimuthal equidistant coords: R_C
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■ Defn: $\Omega_k := -\frac{c^2 k}{a^2 H^2}$ $\Rightarrow k = -\frac{\Omega_{k0} H_0^2}{c^2}$

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■ $\Omega_{\text{tot}0} > 1$

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■ $\Omega_{\text{tot}0} > 1$ *spherical*

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■ $\Omega_{\text{tot}0} > 1$ *spherical* R_C real

FLRW curvature constant

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■ $\Omega_{\text{tot}0} > 1$ *spherical* R_C real

■ $\Omega_{\text{tot}0} = 1$ *flat*

FLRW curvature constant

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- ◆ orthographic coords: k

■ orthographic: $1 - kr^2 = 0$ coord singularity at equator

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■ $\Omega_{\text{tot}0} > 1$ *spherical* R_C real

■ $\Omega_{\text{tot}0} = 1$ *flat* R_C undefined

FLRW curvature constant

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■ $\Omega_{\text{tot}0} = 1$ *flat* R_C undefined

■ $\Omega_{\text{tot}0} < 1$

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■ $\Omega_{\text{tot}0} > 1$ *spherical* R_C real

■ $\Omega_{\text{tot}0} = 1$ *flat* R_C undefined

■ $\Omega_{\text{tot}0} < 1$ *hyperbolic*

FLRW curvature constant

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■ $\Omega_{\text{tot}0} > 1$ *spherical* R_C real

■ $\Omega_{\text{tot}0} = 1$ *flat* R_C undefined

■ $\Omega_{\text{tot}0} < 1$ *hyperbolic* R_C imaginary (or use $|R_C|$)

Einstein's free parameter: Λ

- Einstein: prevent expansion/contraction via Λ

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- MAXIMA: calculate G and $G - g\Lambda = 8\pi T$ and simplify:
<http://cosmo.torun.pl/Cosmo/FLRWEquationsGR>

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- MAXIMA: calculate G and $G - g\Lambda = 8\pi T$ and simplify:
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- *hint:* mixed index form of g is easy

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Friedmann Eqn ($\Lambda \neq 0$):

$$\frac{c^2 k}{a^2} + \frac{\dot{a}^2}{a^2} = \frac{8\pi G \rho}{3} + \frac{c^2 \Lambda}{3}$$

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Defn: $\Omega_\Lambda := \frac{c^2 \Lambda}{3 H^2}$

acceleration Eqn ($\Lambda \neq 0$):

$$\frac{\ddot{a}}{a} = -\frac{H^2}{2} \frac{\rho}{\rho_{\text{crit}}} + \Omega_\Lambda H^2$$

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acceleration Eqn ($\Lambda \neq 0$):

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Defn: $q := -\frac{\ddot{a}a}{\dot{a}^2}$

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Defn: “dust solution”: $p(t) = 0$

Defn: $\Omega_\Lambda := \frac{c^2 \Lambda}{3 H^2}$

Defn: $q := -\frac{\ddot{a}a}{\dot{a}^2}$ “deceleration parameter”

acceleration Eqn ($\Lambda \neq 0$):

$$q = \frac{\Omega_m}{2} - \Omega_\Lambda$$

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- $$q = \frac{\Omega_m}{2} - \Omega_\Lambda$$
 acceleration equation

- if $\Lambda = 0$ and $\Omega_m > 0$ then $\frac{\ddot{a}}{a} < 0$, i.e. $q > 0$

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- $$q = \frac{\Omega_m}{2} - \Omega_\Lambda$$
 acceleration equation

- if $\Omega_\Lambda > \Omega_m/2$ then $\frac{\ddot{a}}{a} > 0$, i.e. $q < 0$ acceleration

distances in FLRW cosmology

- azimuthal equidistant r :

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- $\underline{\Omega_m} = \frac{\rho}{\rho_{\text{crit}}}$

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- $\underline{\Omega_m} = \frac{\rho}{\rho_{\text{crit}}} = \frac{\rho_0 a^{-3}}{\rho_{\text{crit}0} (H^2/H_0^2)}$

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- $\Omega_k = \Omega_{k0} \dot{a}^{-2} H_0^2$
- $\underline{\Omega_\Lambda} = \Omega_{\Lambda0} a^2 \dot{a}^{-2} H_0^2$

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- $$\Rightarrow \dot{a}^2 = H_0^2 (\Omega_{m0} a^{-1} + \Omega_{k0} + \Omega_{\Lambda0} a^2)$$

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$$\Rightarrow \text{when } 0 < z \ll 1, \beta := v/c, \quad r \approx \frac{z c}{H_0}$$

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 r &= \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{\Omega_{m0}a^{-1} + \Omega_{k0} + \Omega_{\Lambda0}a^2}} \\
 &\approx - \int_z^0 \frac{c}{H_0} \frac{1}{\sqrt{1}} dz \\
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EdS: radial comoving distance

$$\blacksquare \quad r = \int_{1/(1+z)}^1 \frac{c}{H_0} \frac{da}{a \sqrt{\Omega_{m0}a^{-1} + \Omega_{k0} + \Omega_{\Lambda0}a^2}}$$

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- high-level frontends (e.g. python) should be easy to write

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want to define d_L such that flux is $F = \frac{L}{4\pi d_L^2}$

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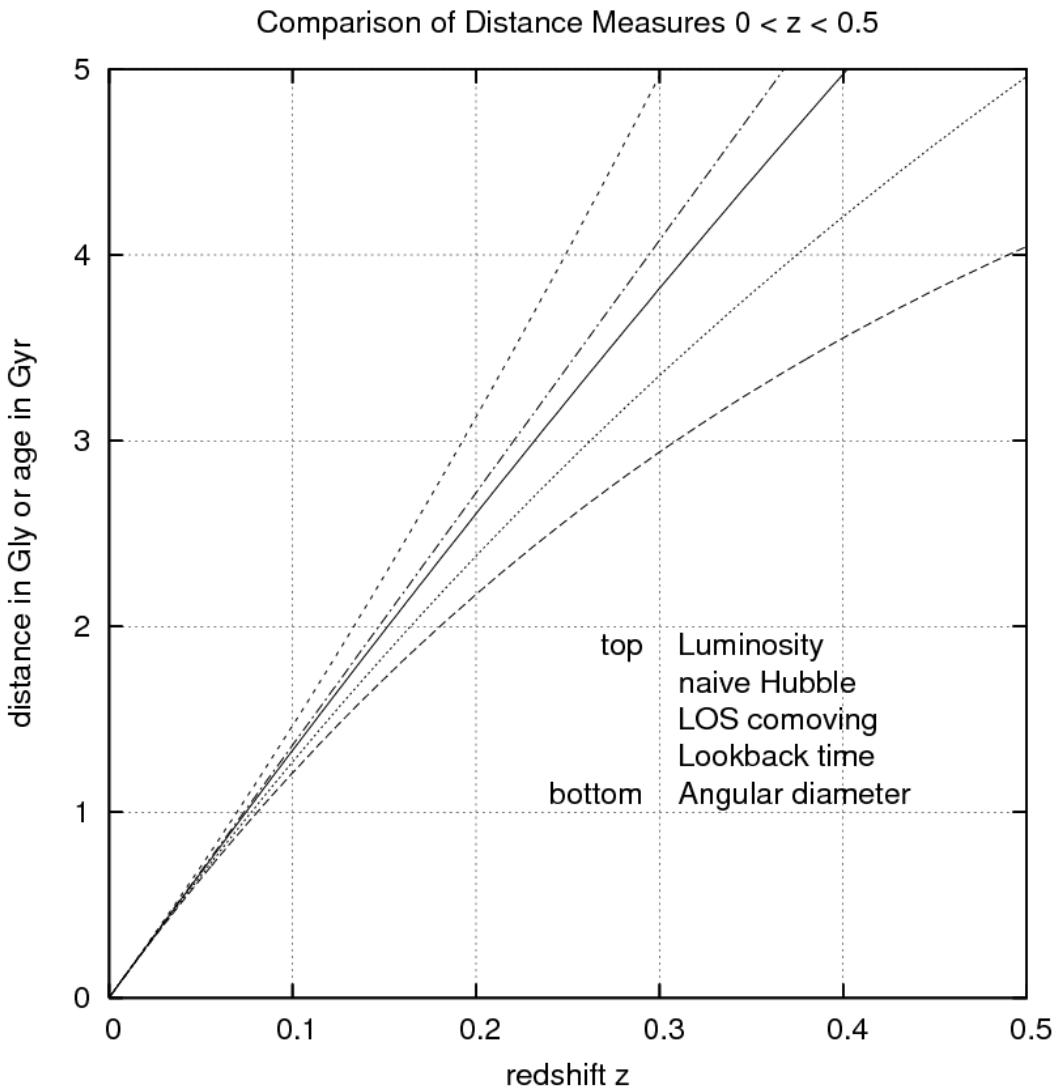
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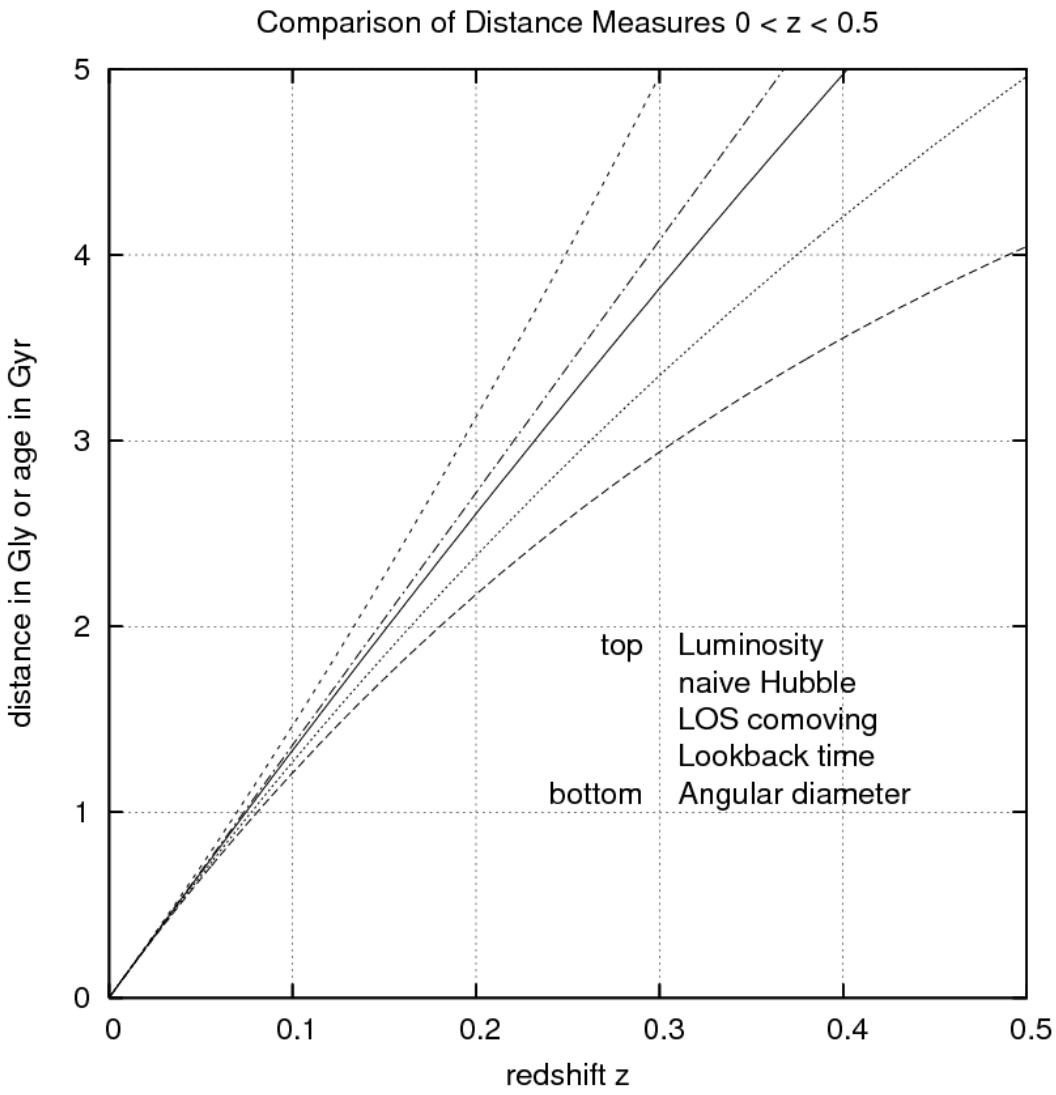
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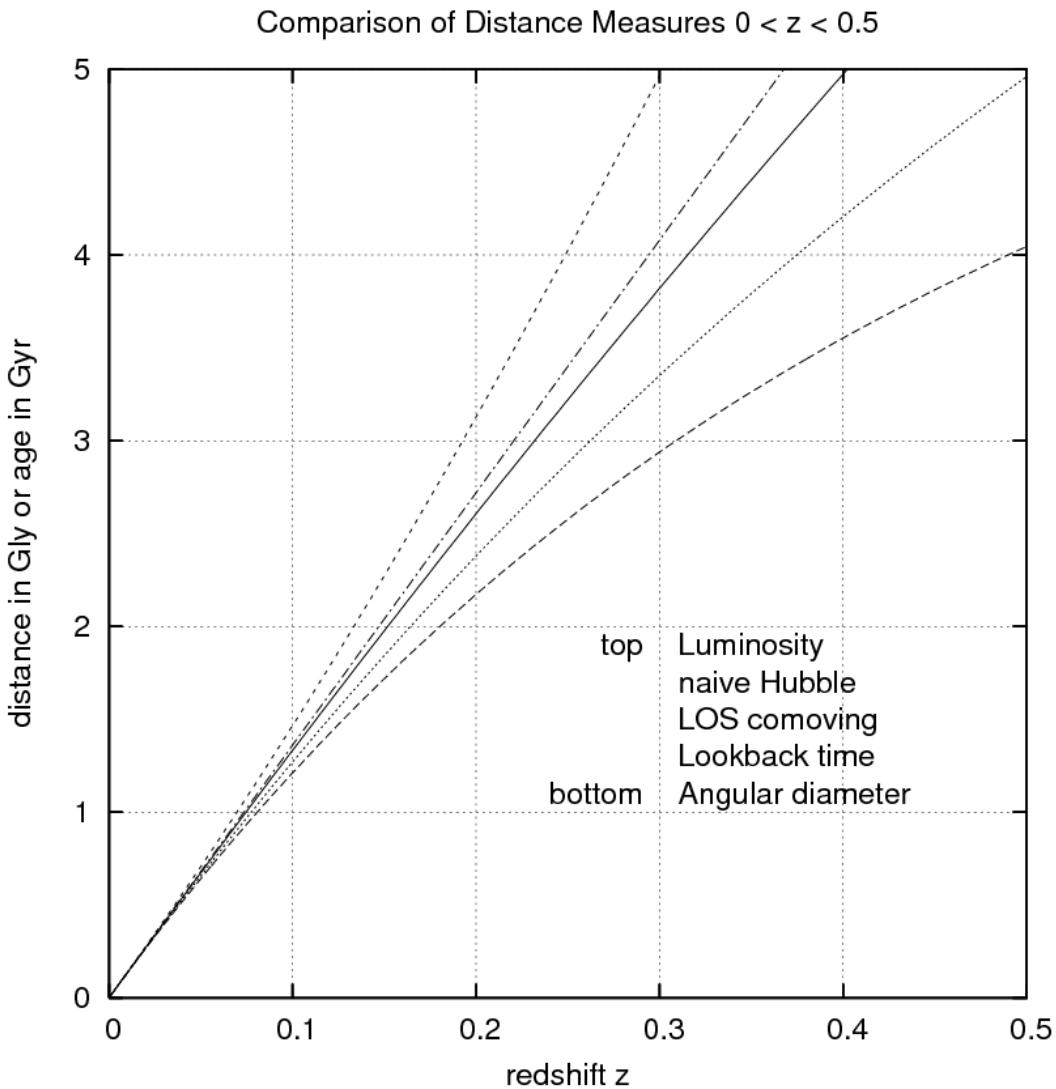


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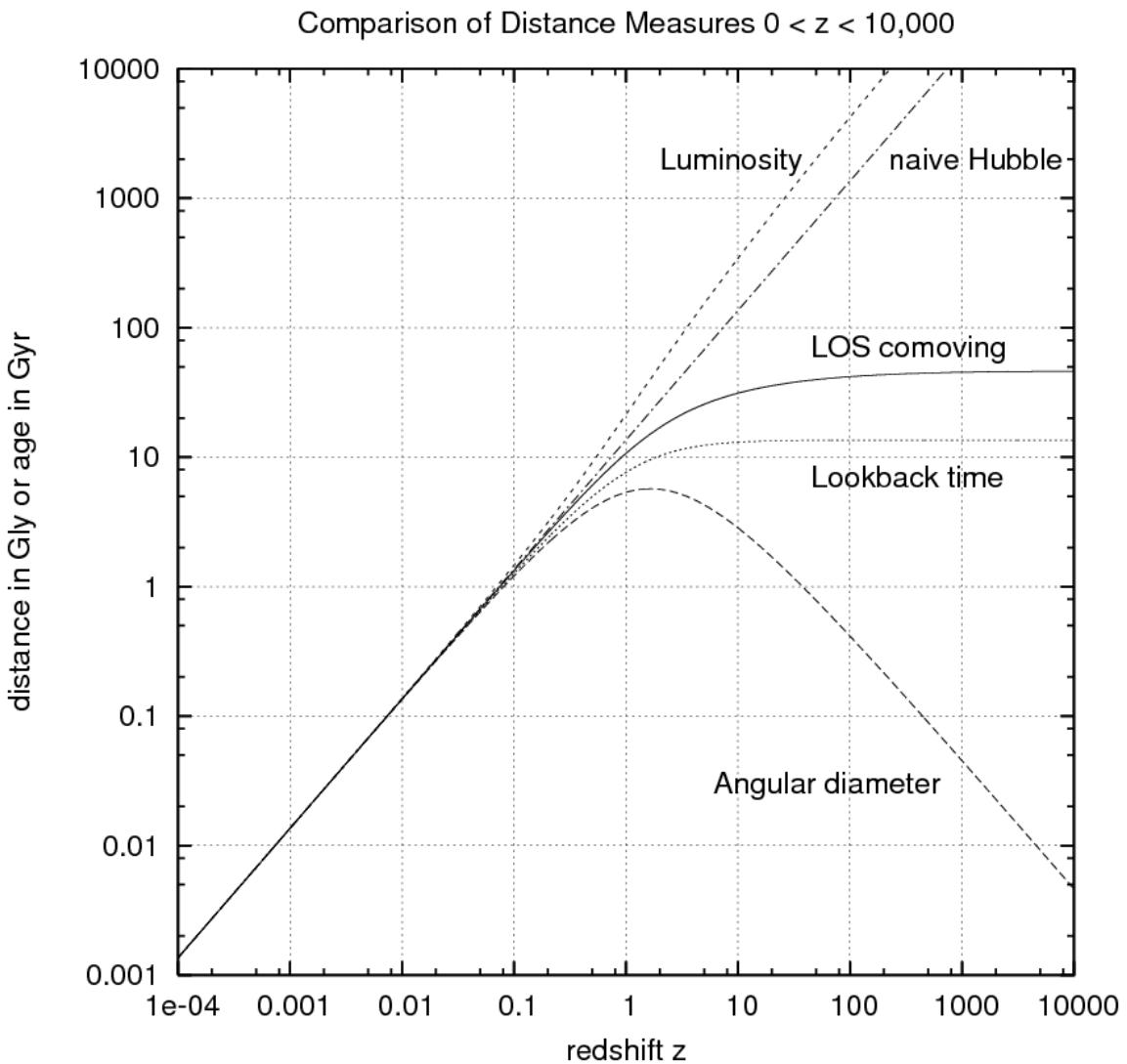
Defn: $h := H_0/100 \text{ km/s/Mpc}$ (without a “0” subscript on h)

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- \Rightarrow no conflict with SR; no conflict with Lorentzian (locally SR) spacetime